

Decomposition of curvature tensor fields in a tachibana first order recurrent space

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Abstract

Takano (1967), have been studied decomposition of curvature tensor in a recurrent space. Sinha and Singh (1970), have been studied and defined decomposition of recurrent curvature tensor field in a Finsler space. Singh and Negi (1985), studied on decomposition of recurrent curvature tensor fields in a Kaehlerian space. Further, Negi and Rawat (1995), studied decomposition in a Kaehlerian recurrent space. Rawat and Silswal (2007), studied decomposition of recurrent curvature fields in a Tachibana space. In the present paper, we considered the decomposition of curvature tensor fields R_{ijk}^h in terms of four non-zero vectors in a Tachibana first order recurrent space. Several theorems also have been established and proved.

Key Words: Recurrent space Curvature tensor, Tachibana Space, First Order.

1.Introduction. An almost Tachibana space is an Almost Hermite space (F_i^h, g_{ij}) , where F_i^h is an Almost

complex structure and g_{ij} is the Hermite metric such that

$$F_{i,j}^h + F_{j,i}^h =, ... (1.1)$$

where the comma (,) followed by an index denotes the operator of covariant differentiation w. r. t. to the Riemannian metric tensor g_{ij} .

In an Almost Tachibana space [10], we have

$$N_{j,i}^h = -4(F_{i,j}^a)F_a^h,$$
 ... (1.2)

where $F_{i,j}^h$ is pure in I and j and N_j^h is Nijenhuis tensor (Yano, 1965), when the Nijenhuis tensor vanishes, the almost Tachibana space is called a Tachibana space and denoted by T_n . A Tachibana space is called Tachibana recurrent space of first order by Lal and Singh,1971, if its curvature tensor field satisfy the condition

$$R_{ijk,a}^h = \lambda_a R_{ijk}^h$$

For some non-zero vector λ_a and called the vector of recurrence.

2. DECOMPOSITION OF RECURRENT CURVATURE TENSOR FIELD R_{ijk}^h

We consider the decomposition of recurrent curvature tensor field R_{ijk}^h in the following form

$$R_{ijk}^h = v^h X_{,i} \phi_j \psi_k , \qquad \dots (2.1)$$

Where the vectors v^h , $X_{,i}$, ϕ_j and ψ_k are such that $\lambda_h v^h = 1 \ . \tag{2.2}$

Theorem 2.1 Under the decomposition (2.1), the Bianchi identity for R_{ijk}^h takes the forms

$$X_{,i}\phi_{j}\psi_{k} + X_{,j}\phi_{k}\psi_{i} + X_{,k}\phi_{i}\psi_{j} = 0,$$
 ... (2.3)

and

$$\lambda_a \phi_j \psi_k + \lambda_j \phi_k \psi_a + \lambda_k \phi_a \psi_j = 0. \qquad \dots (2.4)$$

Proof: From (2.1), we have

$$X_i \phi_i \psi_k + X_j \phi_k \psi_i + X_k \phi_i \psi_i = 0 \qquad \dots (2.5)$$

(since $v^h \neq 0$)

From (1.3) and (2.1), we have

$$v^h X_{,i} [\lambda_a \phi_j \psi_k + \lambda_j \phi_k \psi_a + \lambda_k \phi_a \psi_j] = 0, \qquad \dots (2.6)$$

Multiplying (2.6) by λ_k and using (2.2), we get

$$X_{i}[\lambda_{a}\phi_{j}\psi_{k} + \lambda_{j}\phi_{k}\psi_{a} + \lambda_{k}\phi_{a}\psi_{j}] = 0 \qquad \dots (2.7)$$

(since $X_{i} \neq 0$)

Or
$$\lambda_a \phi_j \psi_k + \lambda_j \phi_k \psi_a + \lambda_k \phi_a \psi_j = 0$$

This completes the proof of the theorem.

Theorem 2.2 Under the decomposition (2.1), the tensor fields R_{ijk}^h , R_{ij} and vectors $X_{,i}$, ϕ_j and ψ_k satisfies the relations

$$\lambda_a R_{ijk}^a = \lambda_i R_{jk} - \lambda_j R_{ik} = X_{,i} \phi_j \psi_k \qquad \dots (2.8)$$

Proof: With the help of Bianchi Identity, equation (1.3) yields

$$\lambda_a R_{ijk}^a = \lambda_i R_{jk} - \lambda_j R_{ik} \qquad \dots (2.9)$$

Multiplying (2.1) by λ_h and using relation (2.2), we get

$$\lambda_h R_{ijk}^a = X_{,i} \phi_j \psi_k \qquad \dots (2.10)$$

From equation (2.9) and (2.10) we get the required relation (2.8).

Theorem 2.3 Under the decomposition (2.1), the quantities λ_a and v^h behave like the recurrent vectors, the recurrent form of these quantities are given by

$$\lambda_{a,m} = \mu_m \lambda_a \qquad \dots (2.11)$$

And
$$v_m^h = -\mu_m v^h$$
 ... (2.12)

Proof: Differentiating (2.9) covariantly w. r. to x^m and using (2.1) and (2.8), we have

$$\lambda_{a,m} v^a X_{,i} \phi_i \psi_k = \lambda_{i,m} R_{ik} - \lambda_{i,m} R_{ik} \qquad \dots (2.13)$$

Multiplying (2,13) by λ_a and using (2.1) and (2.9), we obtain

$$\lambda_{a,m}(\lambda_i R_{ik} - \lambda_j R_{ik}) = \lambda_a(\lambda_{i,m} R_{ik} - \lambda_{i,m} R_{ik}) \qquad \dots (2.14)$$

Now, multiplying equation (2.14) by λ_h , we get

$$\lambda_{a,m} (\lambda_i R_{jk} - \lambda_j R_{ik}) \lambda_h = \lambda_a \lambda_h (\lambda_{i,m} R_{jk} - \lambda_{j,m} R_{ik}) \qquad \dots (2.15)$$

Since the expression on the R.H.S. of the above equation is symmetric in a and h, therefore

$$\lambda_{a\,m}\lambda_{h} = \lambda_{h,m}\lambda_{a} \qquad \dots (2.16)$$

Provided that

$$\lambda_i R_{jk} - \lambda_j R_{ik} \neq 0$$

The vector field λ_a being non-zero, we can have a proportion vector μ_m such that

$$\lambda_{a,m} = \mu_m \lambda_a \qquad \dots (2.17)$$

Further, differentiating the equation (2.2) w. r. to x^m and using relation (2.11), we get

$$v_m^h = -\mu_m v^h \text{ (since } \lambda_h \neq 0) \qquad \dots (2.18)$$

This proves the theorem

Theorem 2.4 Under the decomposition (2.1), the vectors $X_{,i}$, ϕ_j and ψ_k satisfies the relations

$$(\lambda_h + \mu_m) X_{,i} \phi_j \psi_k = X_{,im} \phi_j \psi_k + X_{,i} \phi_{j,m} \psi_k + X_{,i} \phi_j \psi_{k,m} \qquad \dots (2.19)$$

Proof: Differentiating (2.1) covariantly w. r. to x^m and using (1.3), (2.1) and (2.12), we obtain the required result (2.19).

Theorem 2.5 Under the decomposition (2.1), the curvature tensor and holomorphically projective curvature tensor are equal if

$$\phi_k \psi_m \{ (X_{,i} \delta_j^h - X_{,j} \delta_i^h) + X_{,l} (F_i^l F_j^h - F_j^l F_i^h) \} + 2F_i^l F_k^h X_{,l} \phi_j \psi_k = 0 \qquad \dots (2.20)$$

Proof: The holomorphically projective curvature tensor field P_{ijk}^h is defined by

$$P_{ijk}^{h} = R_{ijk}^{h} + \frac{1}{(n+2)} \left(R_{ik} \delta_{j}^{h} - R_{jk} \delta_{i}^{h} + S_{ik} F_{j}^{h} - S_{jk} F_{i}^{h} + 2S_{ij} F_{k}^{h} \right), \qquad \dots (2.21)$$

which may be written in the form

$$P_{ijk}^{h} = R_{ijk}^{h} + D_{ijk}^{h}, (2.22)$$

where

$$D_{ijk}^{h} = \frac{1}{(n+2)} \left(R_{ik} \delta_j^h - R_{jk} \delta_i^h + S_{ik} F_j^h - S_{jk} F_i^h + 2 S_{ij} F_k^h \right) \qquad \dots (2.23)$$

Contracting indices h and k in (2.1), we have

$$R_{ij} = v^k X_{,i} \phi_j \psi_k \qquad \dots (2.24)$$

In view of (2.24), we have

$$S_{ij} = F_i^l v^m X_i \phi_i \psi_k \qquad \dots (2.25)$$

Making use of (2.24) and (2.25) in equation (2.23), we have

$$D_{ijk}^{h} = \frac{1}{(n+2)} \left[v^{m} \phi_{k} \psi_{m} \left(X_{,i} \delta_{j}^{h} - X_{,j} \delta_{i}^{h} \right) + v^{m} X_{,l} \phi_{k} \psi_{m} \left(F_{i}^{l} F_{j}^{h} - F_{j}^{l} F_{i}^{h} \right) + 2 F_{i}^{l} F_{k}^{h} X_{,l} m \phi_{j} \psi_{m} \right]$$
... (2.26)

From equation (2.22), it is clear that

$$P_{ijk}^h = R_{ijk}^h$$
 if $D_{ijk}^h = 0$, which in view of equation (2.26) becomes $v^m \phi_k \psi_m (X_{,i} \delta_j^h - X_{,j} \delta_i^h) + v^m X_{,l} \phi_k \psi_m (F_i^l F_j^h - F_j^l F_i^h) + 2 F_i^l F_k^h X_{,l} m \phi_j \psi_m = 0...$ (2.27) Multiplying the above equation by λ_m and using relation (2.2), we obtain the required condition (2.20).

Theorem 2.6 Under the decomposition (2.1), the scalar curvature R, satisfy the relation

$$\lambda_k R = g^{ij} X_i \phi_i \psi_k \qquad \dots (2.28)$$

Proof: Contracting indices h and k in (2.1), we get

$$R_{ij} = v^k X_i \phi_i \psi_k \qquad \dots (2.29)$$

Multiplying (2.29) by g^{ij} on both sides, we get

$$g^{ij}R_{ij} = g^{ij}v^k X_{,i}\phi_j\psi_k \qquad \dots (2.30)$$

Or

$$R = g^{ij}v^k X_i \phi_i \psi_k \quad (R = g^{ij}R_{ii}) \qquad \dots (2.31)$$

Now, multiplying (2.31) by λ_k and using (2.2), we get

$$\lambda_k R = g^{ij} X_{,i} \phi_j \psi_k,$$

which completes the proof of the theorem.

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